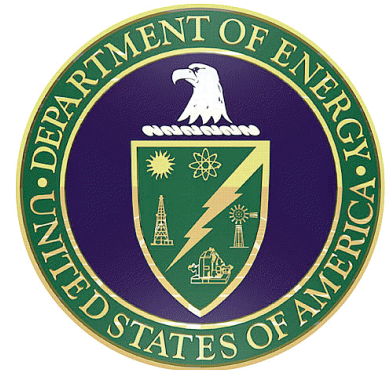


The β - ν Correlation in Laser-Trapped ^{21}Na and the Standard Model

Paul Vetter



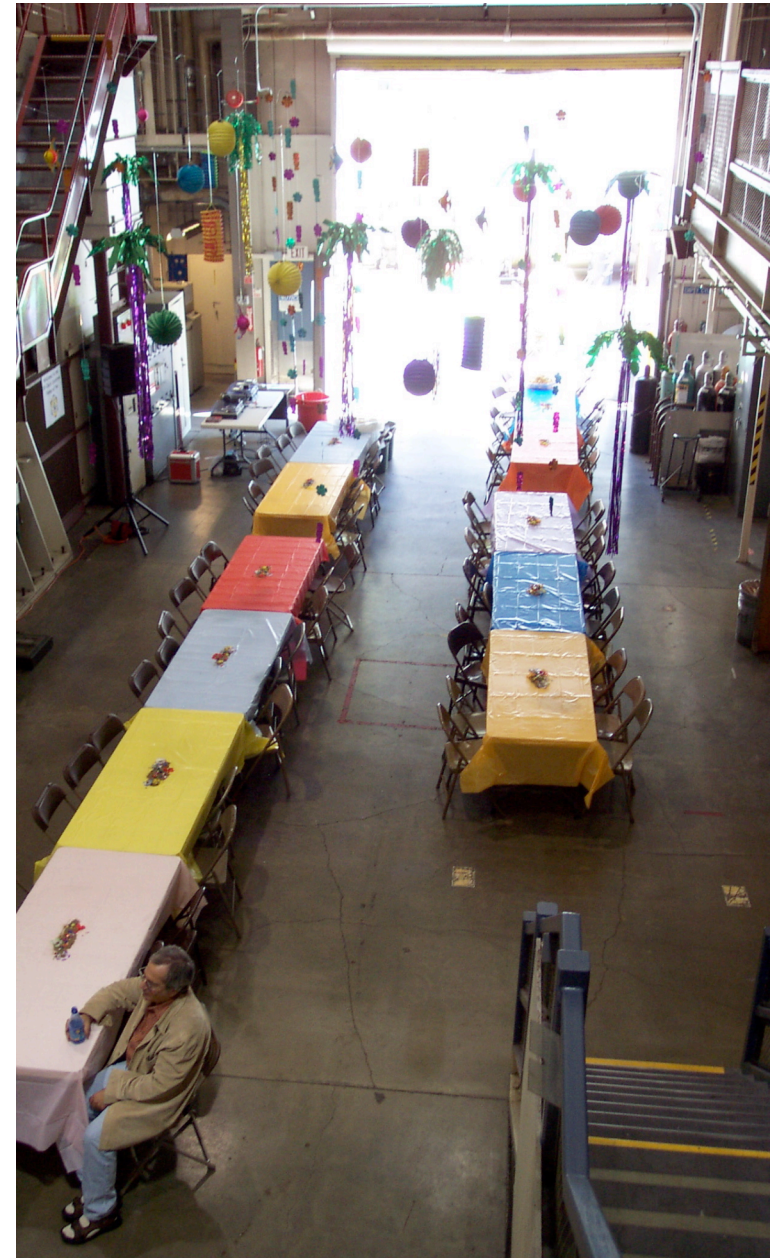
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Eric Oelker

Feb. 25, 2008



The electroweak Lagrangian

$$\begin{aligned}\mathcal{L}_F &= \sum_i \bar{\psi}_i \left(i \not{\partial} - m_i - \frac{gm_i H}{2M_W} \right) \psi_i \\ &- \frac{g}{2\sqrt{2}} \sum_i \bar{\psi}_i \gamma^\mu (1 - \gamma^5) (T^+ W_\mu^+ + T^- W_\mu^-) \psi_i \\ &- e \sum_i q_i \bar{\psi}_i \gamma^\mu \psi_i A_\mu \\ &- \frac{g}{2 \cos \theta_W} \sum_i \bar{\psi}_i \gamma^\mu (g_V^i - g_A^i \gamma^5) \psi_i Z_\mu\end{aligned}$$

$$\mathcal{L}_F = \frac{g}{2\sqrt{2}} \sum_i \bar{\psi}_i \gamma^\mu (1 - \gamma^5) (T^+ W_\mu^+ + T^- W_\mu^-) \psi_i$$

Fermions come in doublets

$$\psi_i = \begin{pmatrix} \nu_i \\ \ell_i^- \end{pmatrix} \qquad \begin{pmatrix} u_i \\ d'_i \end{pmatrix}$$

$$d'_i = \sum_j V_{ij} d_j$$

Where quark mixing is a unitary transform like this:

$$V_{ij} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

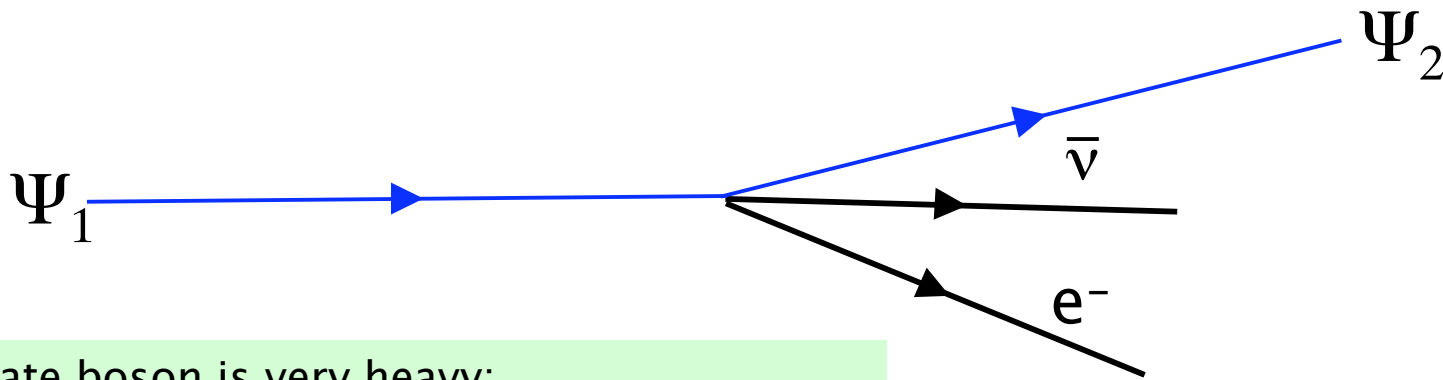
Which can be written as:

$$V_{ij} = \begin{pmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3 (\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3 (1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} \\ + O(\lambda^4)$$

So is this all there is?

$$\mathcal{L}_F = \frac{g}{2\sqrt{2}} \sum_i \bar{\psi}_i \gamma^\mu (1 - \gamma^5) (T^+ W_\mu^+ + T^- W_\mu^-) \psi_i$$

Nuclear beta decay four fermion interaction



Intermediate boson is very heavy:
Point-like interaction of the four fermions' currents

Isospin quantum numbers in the nucleus

$$\begin{aligned}
 H_{\text{int}} = & (\bar{\psi}_2 \psi_1) (C_S \bar{\psi}_e \psi_\nu + C'_S \bar{\psi}_e \gamma_5 \psi_\nu) \\
 & + (\bar{\psi}_2 \gamma_\mu \psi_1) (C_V \bar{\psi}_e \gamma^\mu \psi_\nu + C'_V \bar{\psi}_e \gamma^\mu \gamma_5 \psi_\nu) \\
 & + 1/2 (\bar{\psi}_2 \sigma_{\lambda\mu} \psi_1) (C_T \bar{\psi}_e \sigma^{\lambda\mu} \psi_\nu + C'_T \bar{\psi}_e \sigma^{\lambda\mu} \gamma_5 \psi_\nu) \\
 & - (\bar{\psi}_2 \gamma_\mu \gamma_5 \psi_1) (C_A \bar{\psi}_e \gamma^\mu \gamma_5 \psi_\nu + C'_A \bar{\psi}_e \gamma^\mu \psi_\nu) \\
 & + (\bar{\psi}_2 \gamma_5 \psi_1) (C_P \bar{\psi}_e \gamma_5 \psi_\nu + C'_P \bar{\psi}_e \psi_\nu)
 \end{aligned}$$

Calculate the decay rate

$$\frac{d^4W(\mathbf{p}_e, \mathbf{p}_\nu, \mathbf{I})}{d\Omega_e d\Omega_\nu} = \xi \left\{ 1 + a_{\beta\nu} \frac{\mathbf{p}_e \cdot \mathbf{p}_\nu}{E_e E_\nu} + b_{\text{Fierz}} \frac{m_e}{E_e} + \frac{\langle \mathbf{I} \rangle}{I} \cdot \left[A_{\beta} \frac{\mathbf{p}_e}{E_e} + B_{\nu} \frac{\mathbf{p}_\nu}{E_\nu} + D \frac{\mathbf{p}_e \times \mathbf{p}_\nu}{E_e E_\nu} \right] \right. \\ \left. + c_{\text{align}} \left[\frac{\mathbf{p}_e \cdot \mathbf{p}_\nu}{3E_e E_\nu} - \frac{(\mathbf{p}_e \cdot \hat{i})(\mathbf{p}_\nu \cdot \hat{i})}{E_e E_\nu} \right] \left[\frac{I(I+1) - 3\langle (\mathbf{I} \cdot \hat{i})^2 \rangle}{I(2I+1)} \right] \right\}$$

+ terms with β polarization

Jackson, Treiman, Wyld, *Phys. Rev.* **106**, p. 517 (1957)

Correlation coefficients in terms of four-fermion coupling constants

Jackson, Treiman and Wyld, *Phys. Rev.* **106**, 517 (1957).

$$\xi = |M_F|^2(|C_S|^2 + |C_V|^2 + |C_S'|^2 + |C_V'|^2) + |M_{GT}|^2(|C_T|^2 + |C_A|^2 + |C_T'|^2 + |C_A'|^2),$$

$$a\xi = |M_F|^2(-|C_S|^2 + |C_V|^2 - |C_S'|^2 + |C_V'|^2) + \frac{|M_{GT}|^2}{3}(|C_T|^2 - |C_A|^2 + |C_T'|^2 - |C_A'|^2),$$

$$b\xi = \pm 2 \operatorname{Re}[|M_F|^2(C_S C_V^* + C_S' C_V'^*) + |M_{GT}|^2(C_T C_A^* + C_T' C_A'^*)],$$

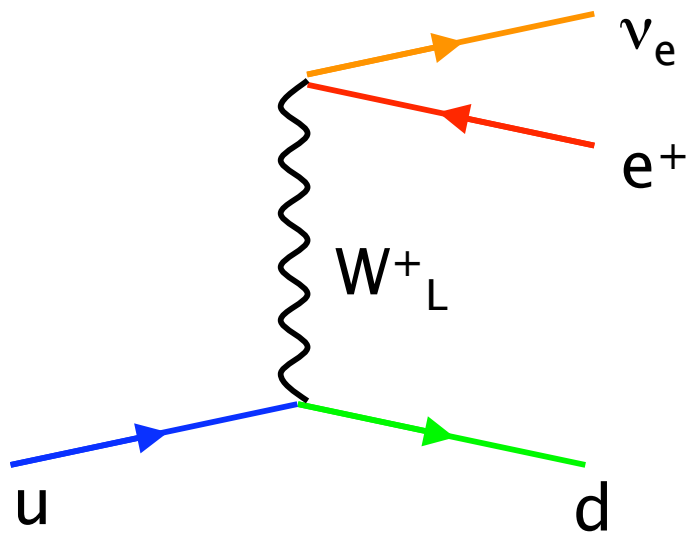
Helicity Projection Formalism

P. Herczeg, *Prog. Nucl. Part. Phys.* **46**, 213 (2001).

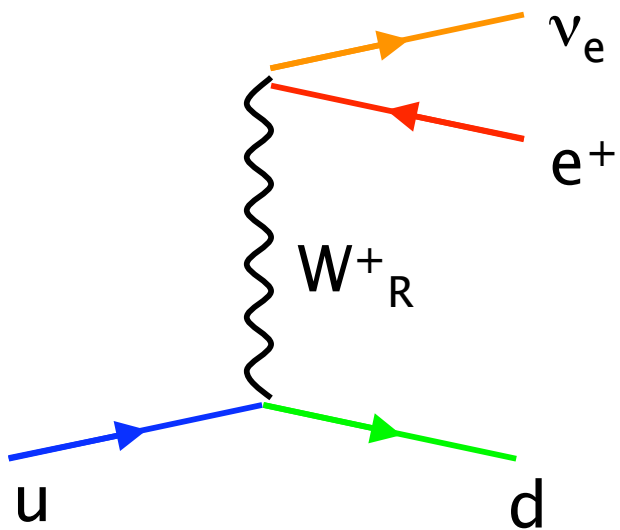
Severijns, Beck and Naviliat-Cuncic, *Rev. Mod. Phys.* **78**, 991 (2006).

Relate C_i 's to g_S a_{LR} A_{LR}

Michel parameter related formalism, more coefficients (flavor dependence)...

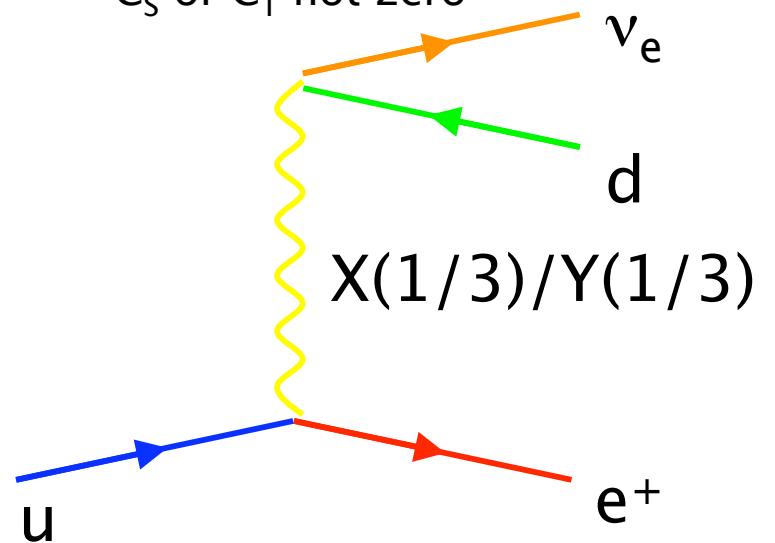


Right handed W boson -- C' s not zero



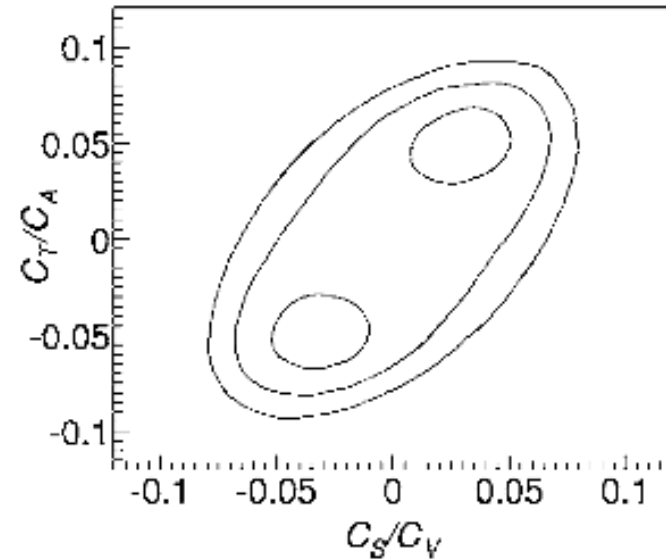
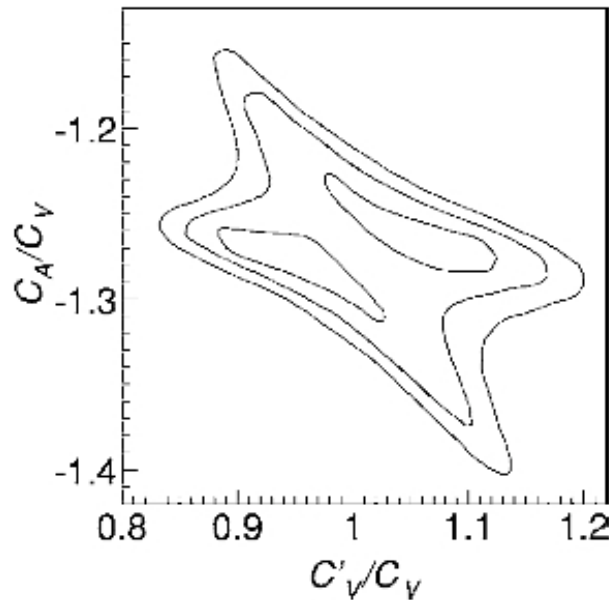
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Scalar or tensor leptoquark
 C_S or C_T not zero



Limits on Non-Standard Model Couplings

N. Severijns, M. Beck, and O. Naviliat-Cuncic,
Rev. Mod. Phys. 78, 991 (2006)



Two possible exclusion plots
Non-StdMod couplings can be ~10% of allowed
Best fit with neutron data suggests new physics

Scalar, tensor couplings can be induced by leptoquarks, left-right symmetry
Supersymmetry, charged Higgs exchange.

P. Herczeg, Prog. Part. Nucl. Phys. 46, 413 (2001).

- Available model space to test is **very large -- many degrees of freedom**
- High energy tests do not span the model space
- Current limits on non-Standard couplings are generally not stringent

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Beta neutrino correlation searches for Scalar, Tensor couplings

$$a_{\beta\nu} = \frac{|M_F|^2(|C_V|^2 + |C'_V|^2 - |C_S|^2 - |C'_S|^2) - \frac{1}{3}|M_{GT}|^2(|C_A|^2 + |C'_A|^2 - |C_T|^2 - |C'_T|^2)}{|M_F|^2(|C_V|^2 + |C'_V|^2 + |C_S|^2 + |C'_S|^2) + |M_{GT}|^2(|C_A|^2 + |C'_A|^2 + |C_T|^2 + |C'_T|^2)}$$

$$M_F = \langle 1 \rangle = \langle \Psi_f | \tau_{\pm} | \Psi_i \rangle$$

$$M_{GT} = \langle \sigma \rangle = \langle \Psi_f | \vec{\sigma} \tau_{\pm} | \Psi_i \rangle$$

Calculate **a** in the Standard Model as:

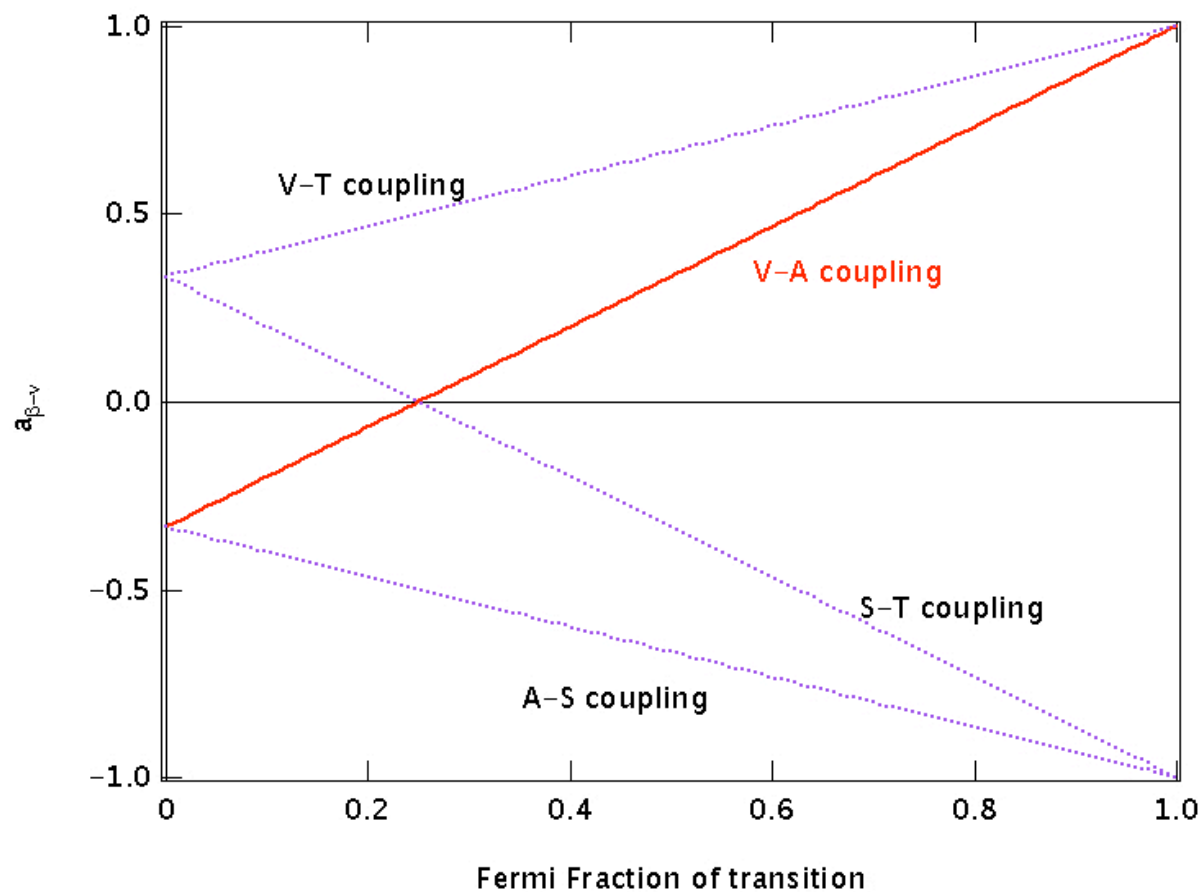
$$a_{\beta\nu} = \frac{1 - \lambda^2/3}{1 + \lambda^2}$$

Where λ is the ratio of axial vector to vector (G-T to Fermi) in the transition

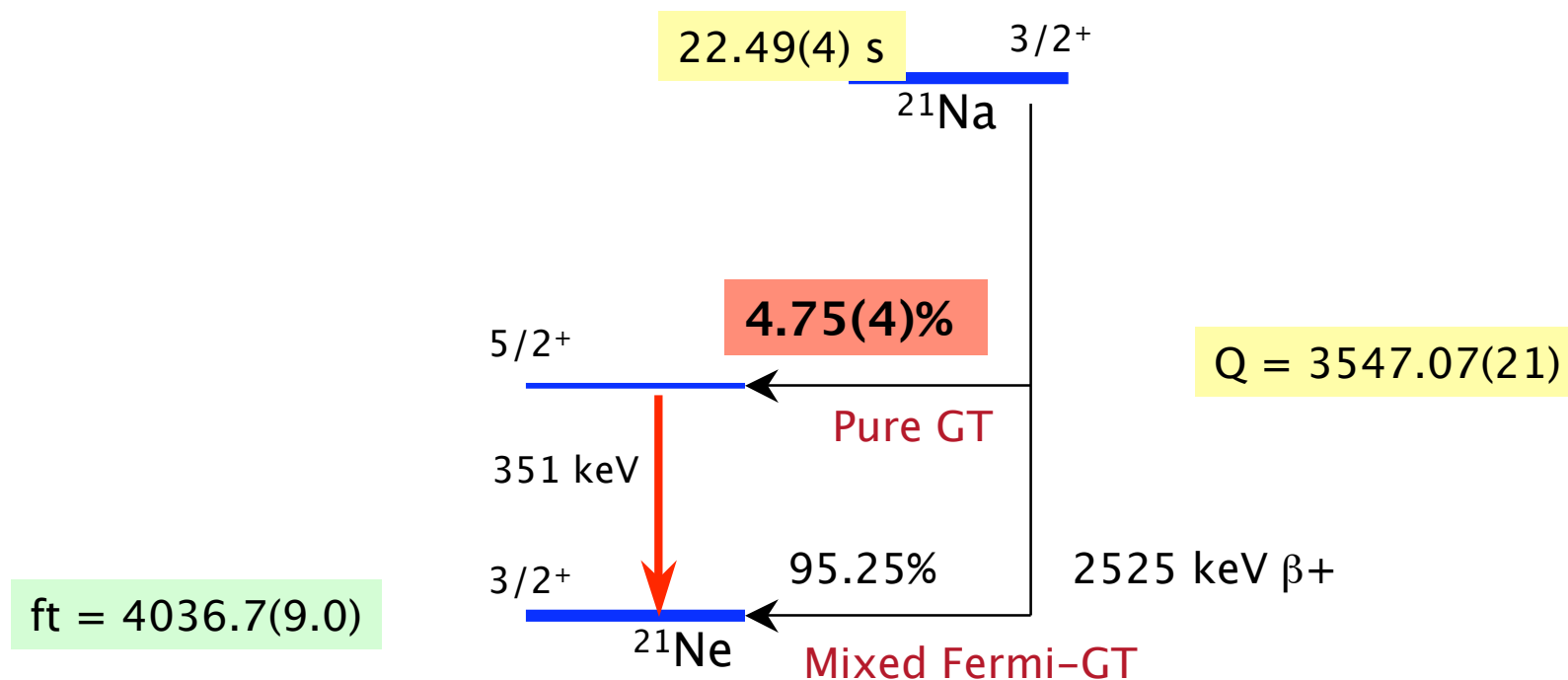
$$\lambda = \frac{C_A \langle \sigma \rangle}{C_V \langle 1 \rangle}$$

Beta Neutrino Correlation Measurements

$a_{\beta\nu}$ plotted as a function of the Fermi fraction of the transition



Decay scheme of ^{21}Na

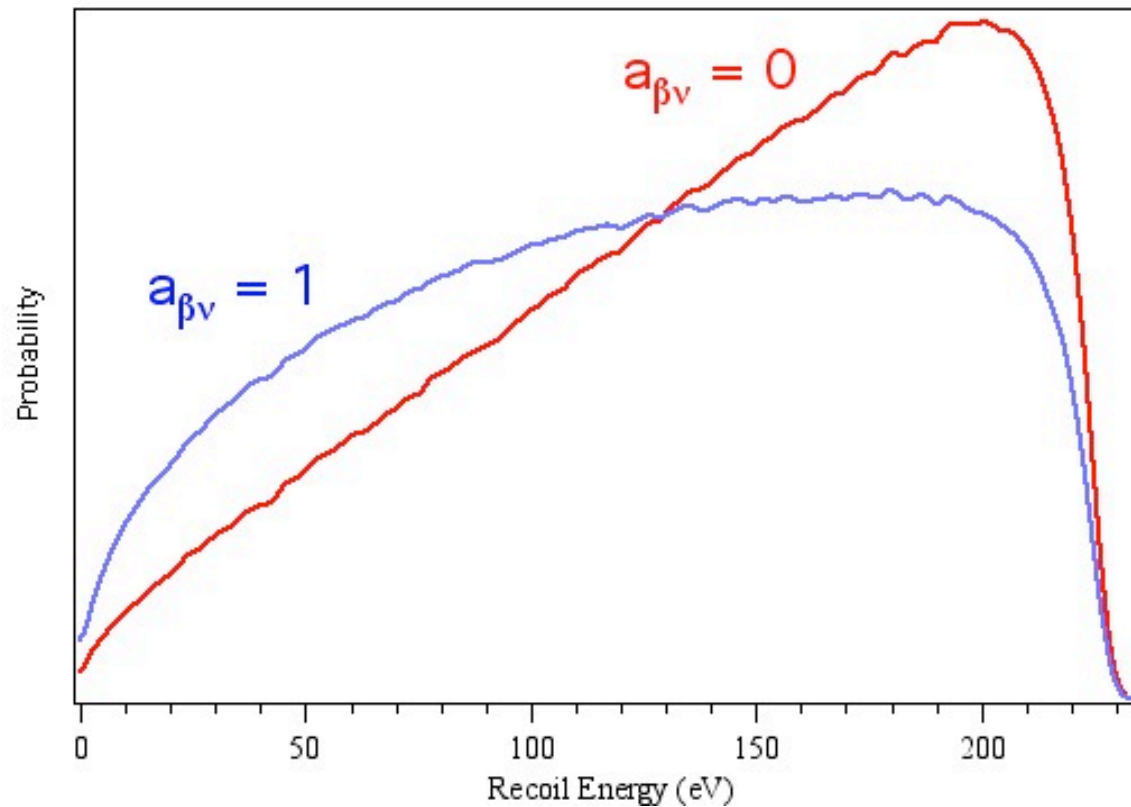


$^{21}\text{Na } 3/2^+ \rightarrow ^{21}\text{Ne } 3/2^+$:

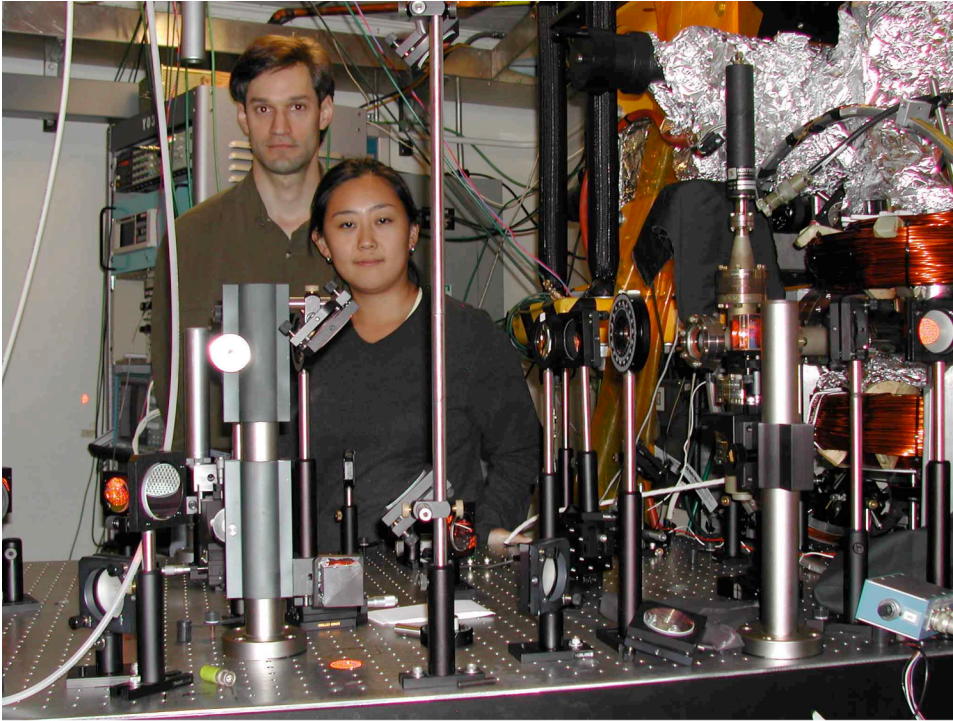
$\langle 1 \rangle = 1$ Fermi matrix element -- isodoublet mirror transition

$C_A \langle \sigma \rangle = 0.703(3)$ Gamow -Teller matrix element (calculated)

Recoil energy spectrum for ^{21}Na



Apparatus



Neutral atom trap for ^{21}Na

Measure total energy/momentum spectrum of recoil nuclei

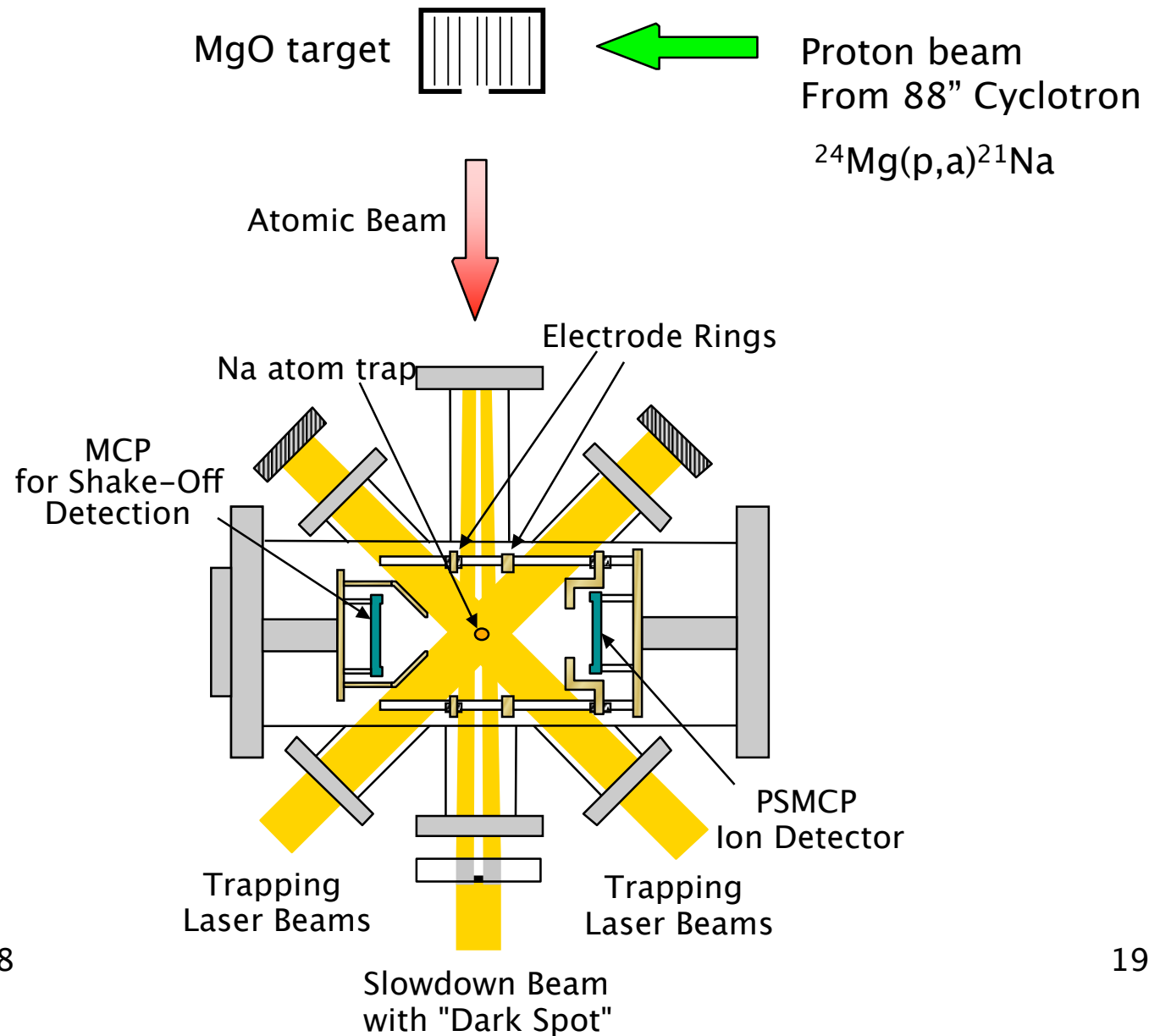
Time-of-flight momentum spectrometer

T=0 start provided by shake-off electrons
High efficiency, run with low trap density

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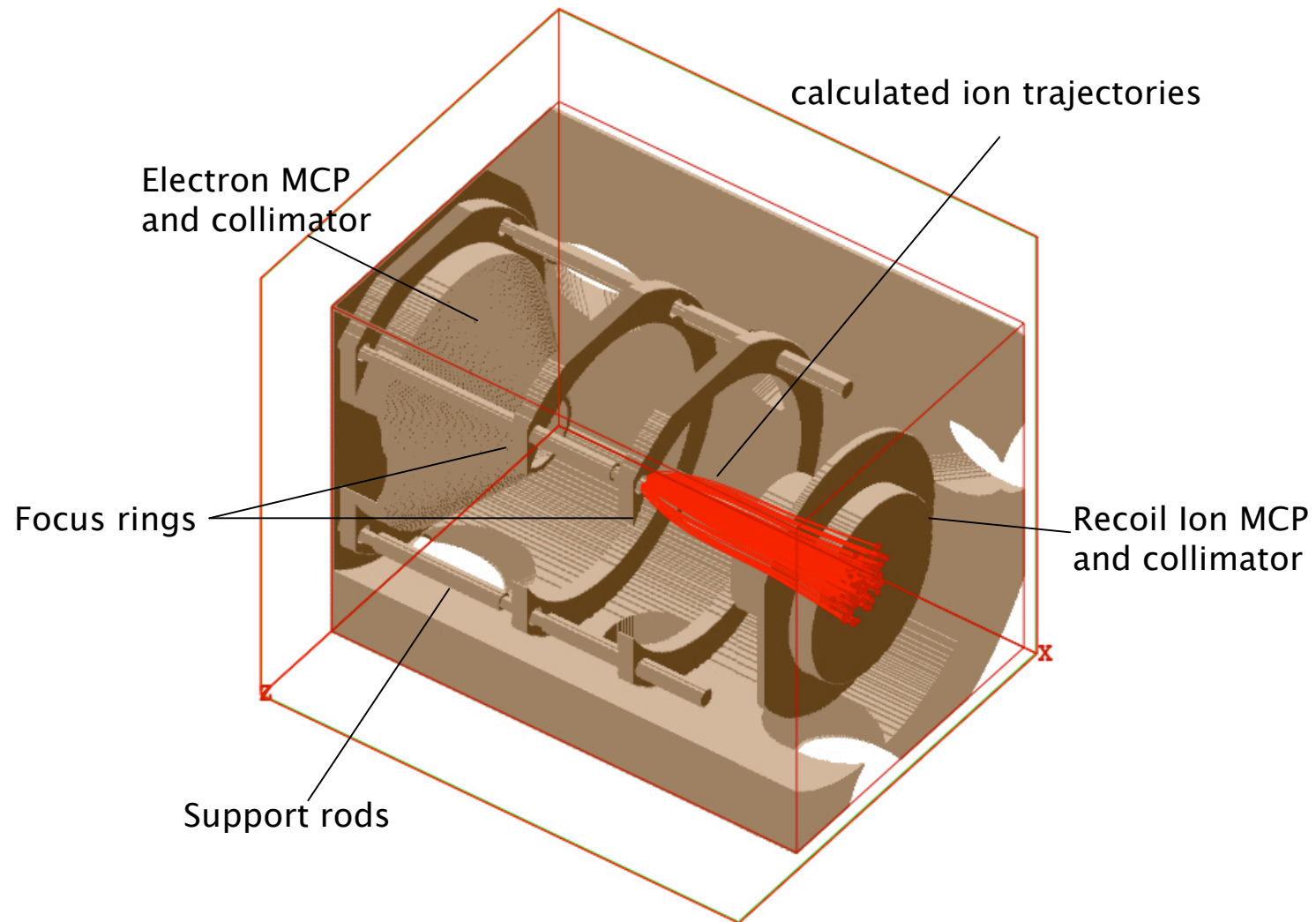
Apparatus schematic



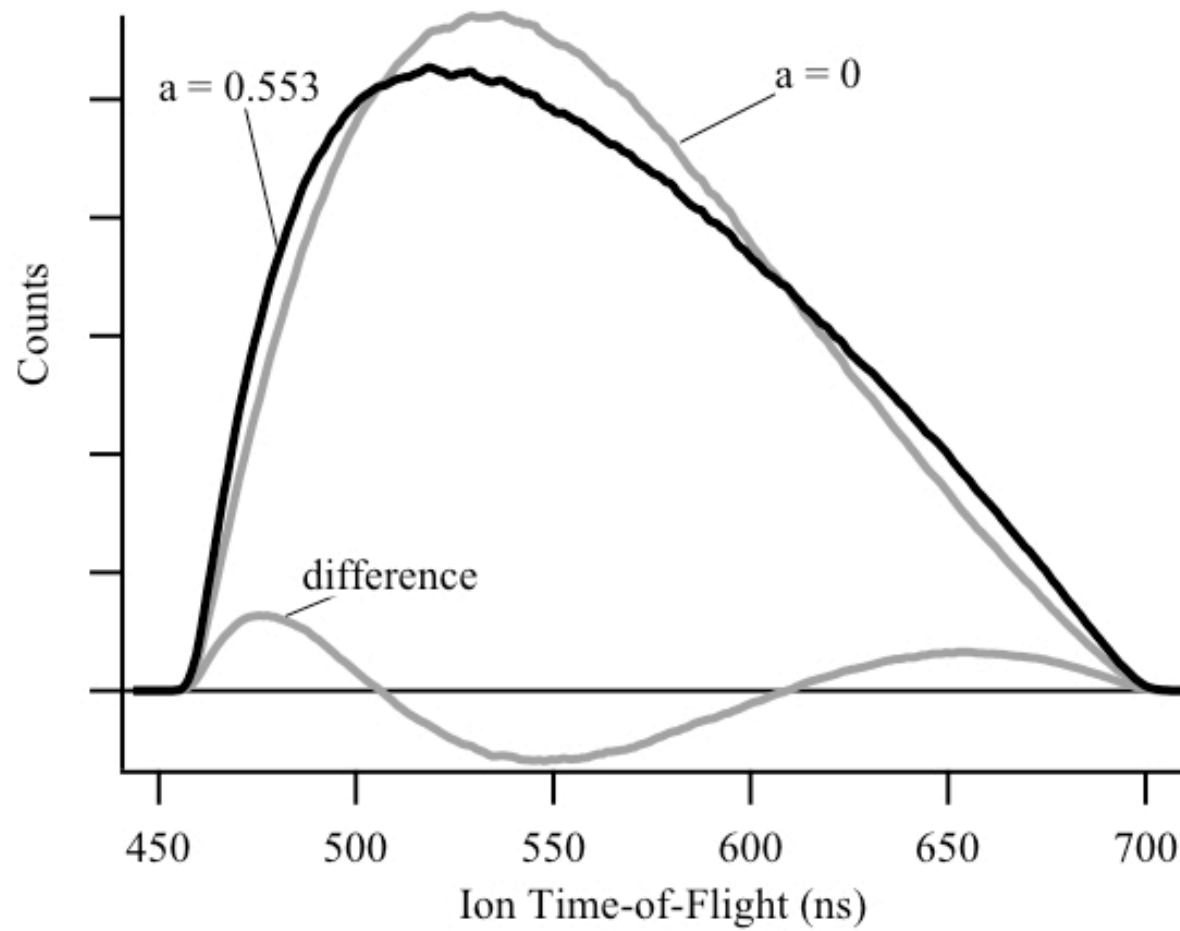
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Simlon

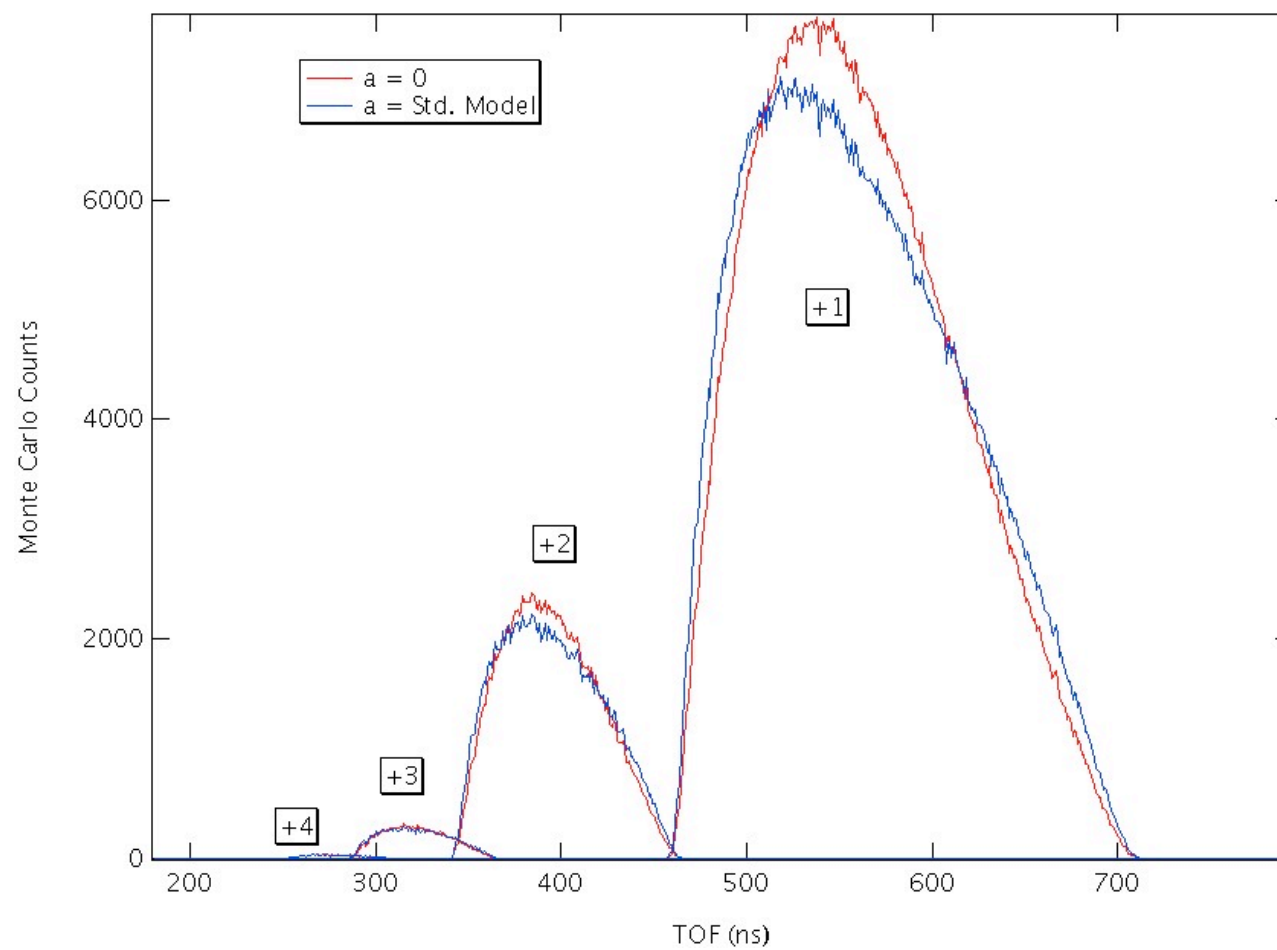
Cutaway view of the simulation volume inside the trapping vacuum chamber.



Calculated Time-of-Flight spectrum (Monte-Carlo)

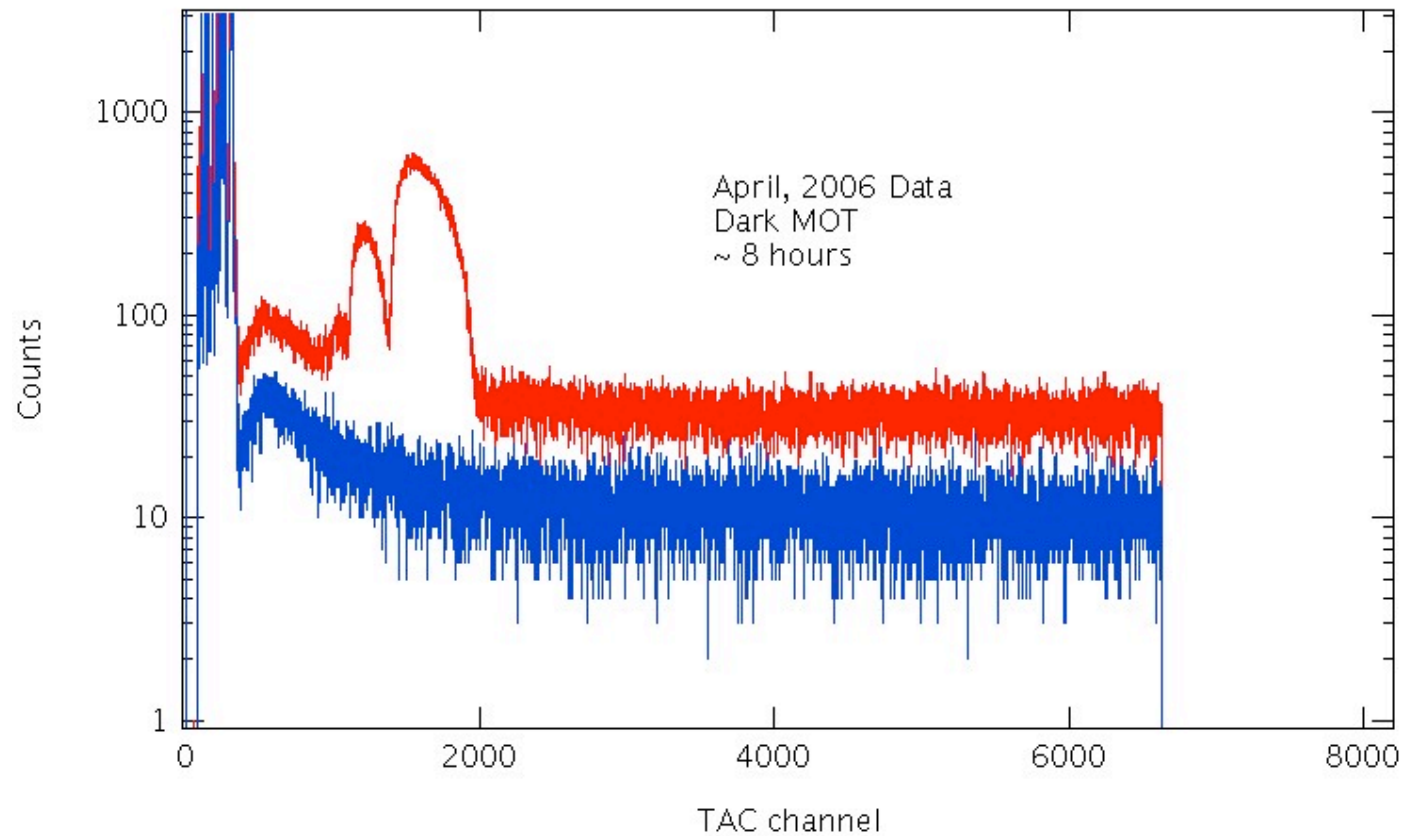


Several possible charge states of ^{21}Ne

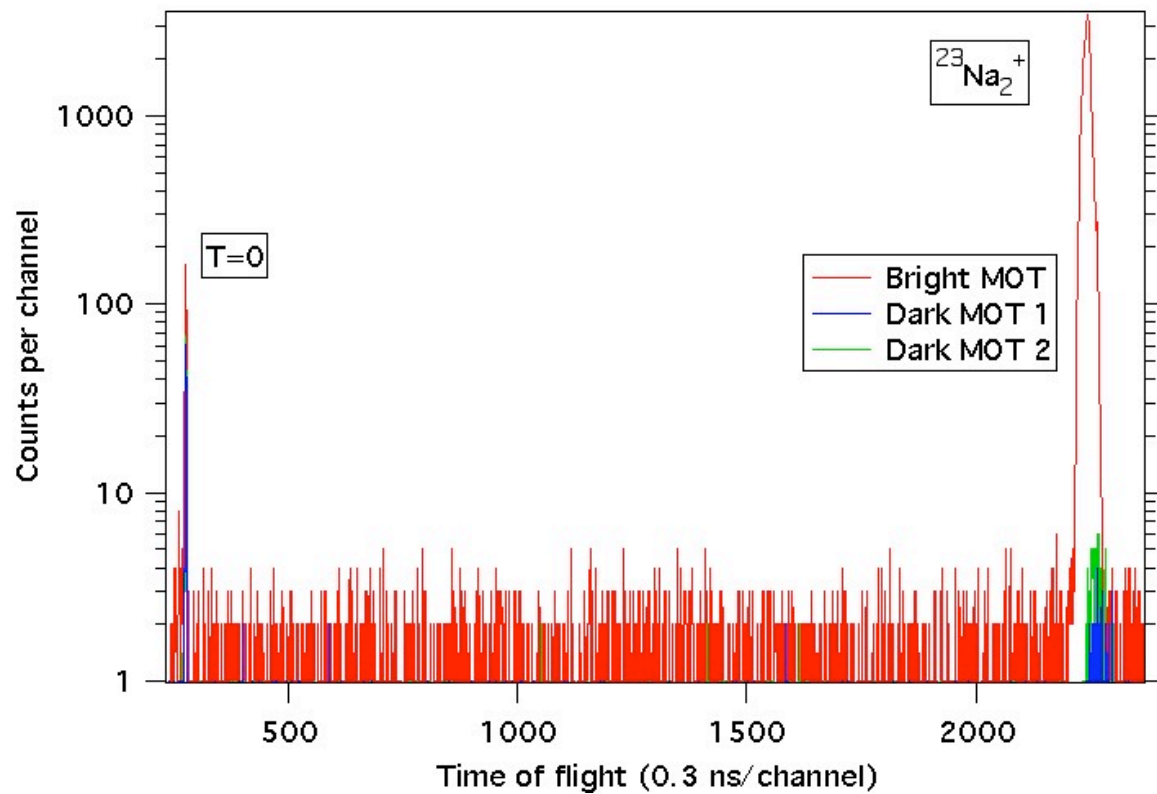


Time of Flight Data

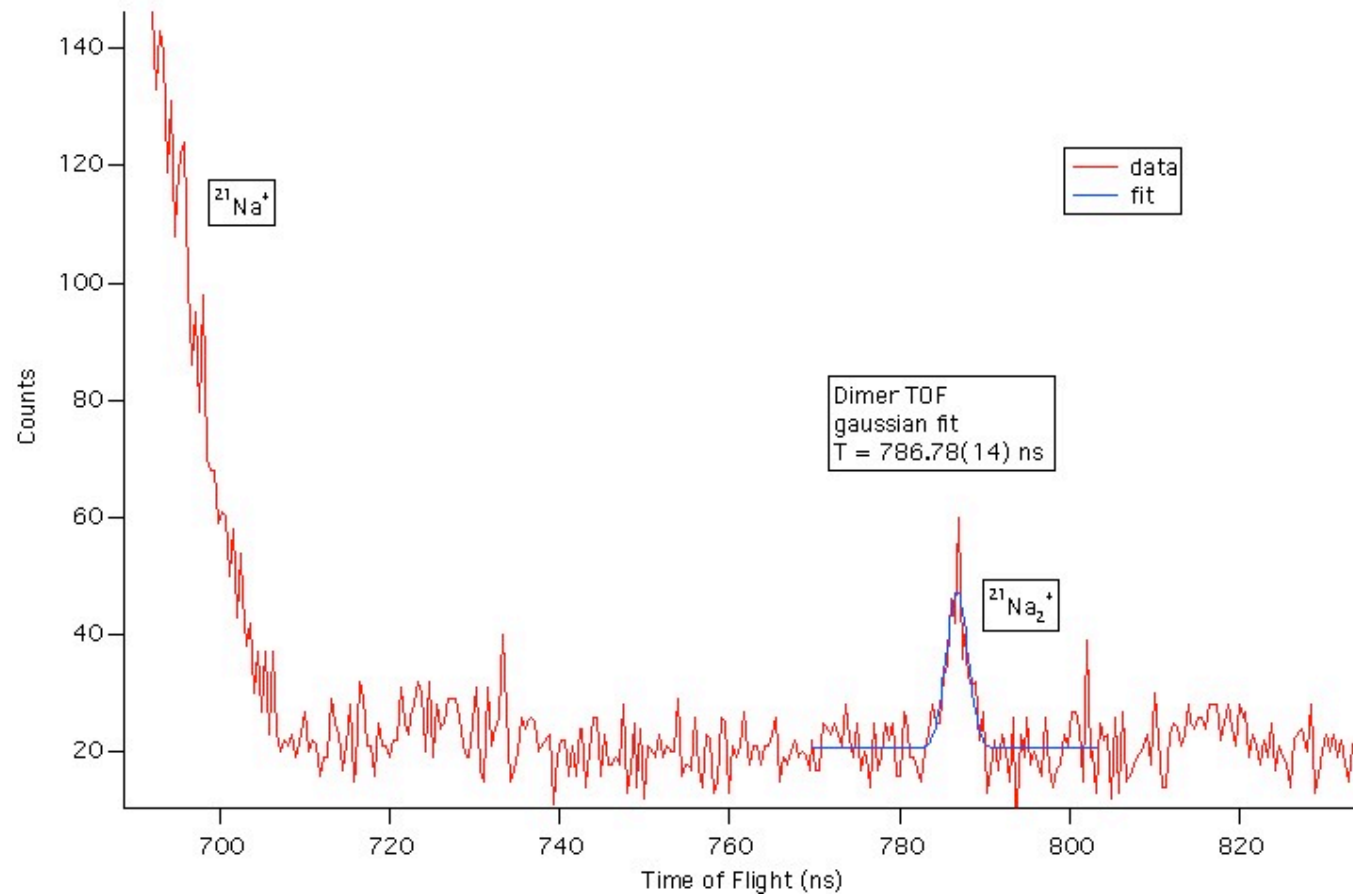
Representative Data



Surprise: Autoionized dimers of $^{23}\text{Na}_2^+$



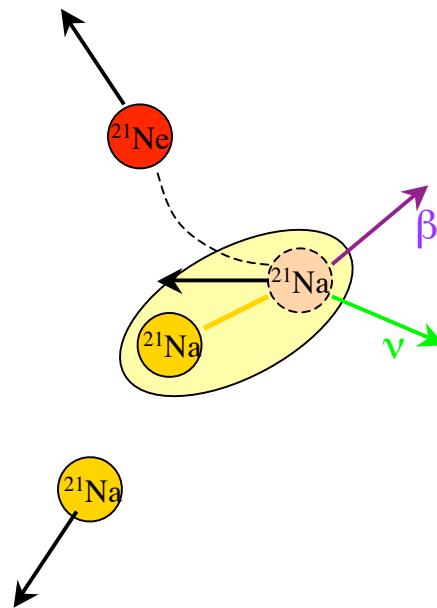
Radioactive Molecules



Beta decay in $^{21}\text{Na}_2$

$^{21}\text{Na}_2$ Formation rates in MOT are high

- Recoil ion scatters off interatomic potential at close range, changes momentum
- Initial momentum, energy distribution of ^{21}Ne is lost
- Depending on scattering parameters (b , θ) and $^{21}\text{Ne}^{+n} - ^{21}\text{Na}$ potential

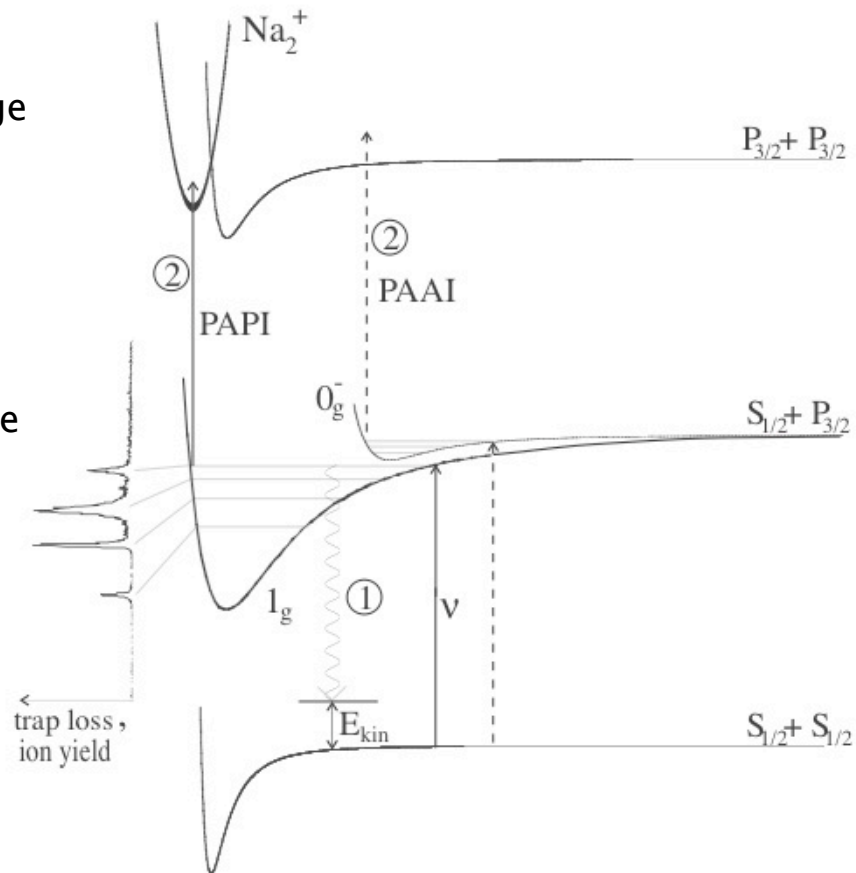


Photoassociation and Autoionization in Sodium

PAPI = Photoassociative Photoionization
Depends mostly on formation of short-range bound states

PAAI = Photoassociative Autoionization
Long-range bound states, not possible in single-color red-detuned experiment

PAPI dominates in Na_2 , implying short-range excited states are populated.
These would be the most significant path to populate cold, bound, ground states

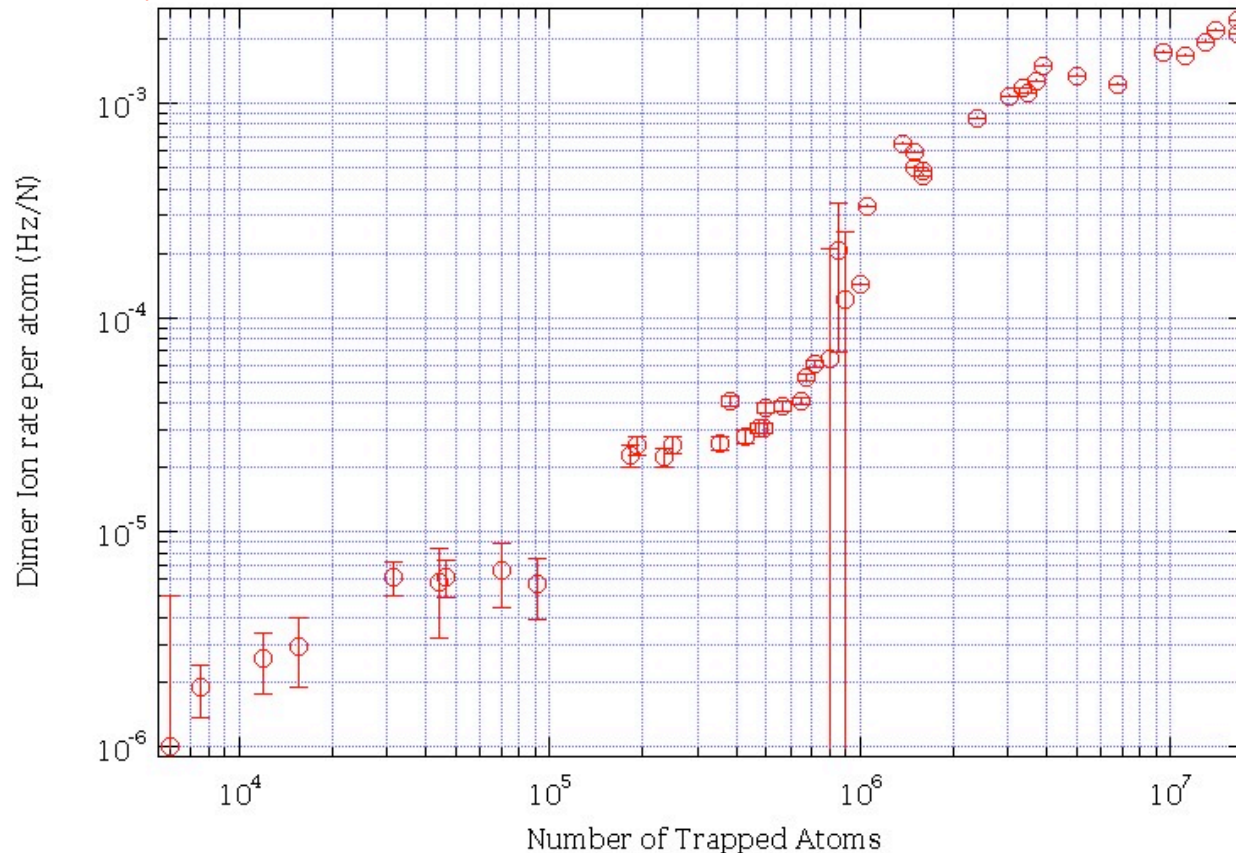


Do small traps reduce the number of molecules?

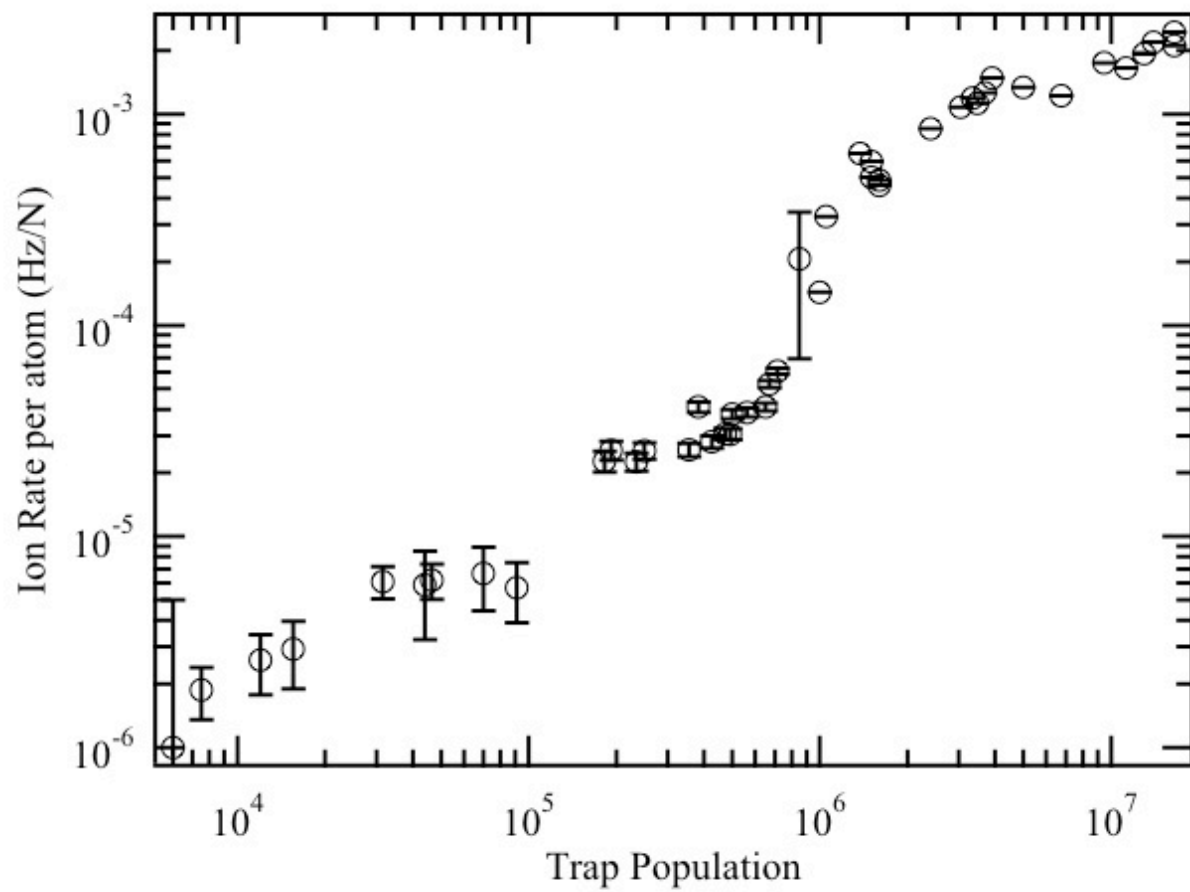
Number of ionized dimer $^{23}\text{Na}_2^+$ molecules detected with MCP's vs. atom trap population

Number of molecules relative to the trap population ($N_{\text{mol}}/N_{\text{atom}}$)

Perturbation on $a_{\beta v} \propto$ number of counts from molecular decays relative to counts from atomic decays



The number of dimers relative to the atomic population falls steeply at small trap sizes.

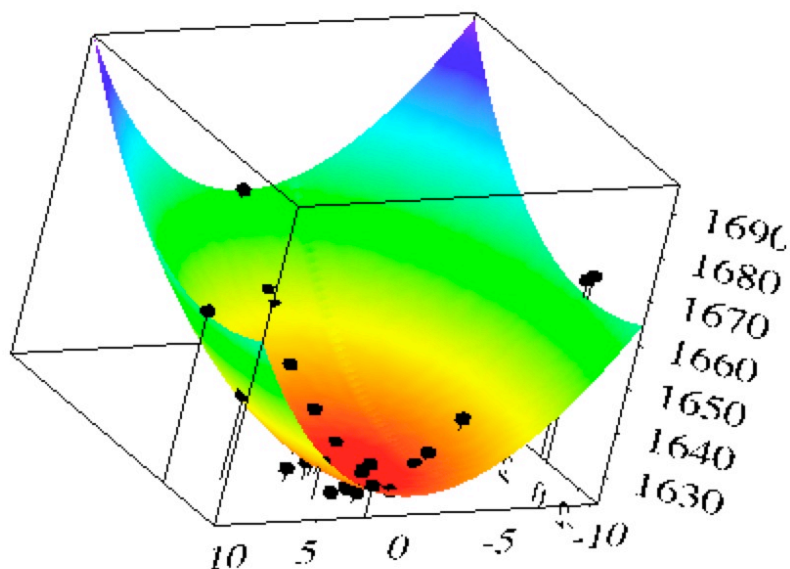


Sometimes molecules are useful Determining the Trap Position

Measure the stable dimer ions' TOF and positions on the MCP.

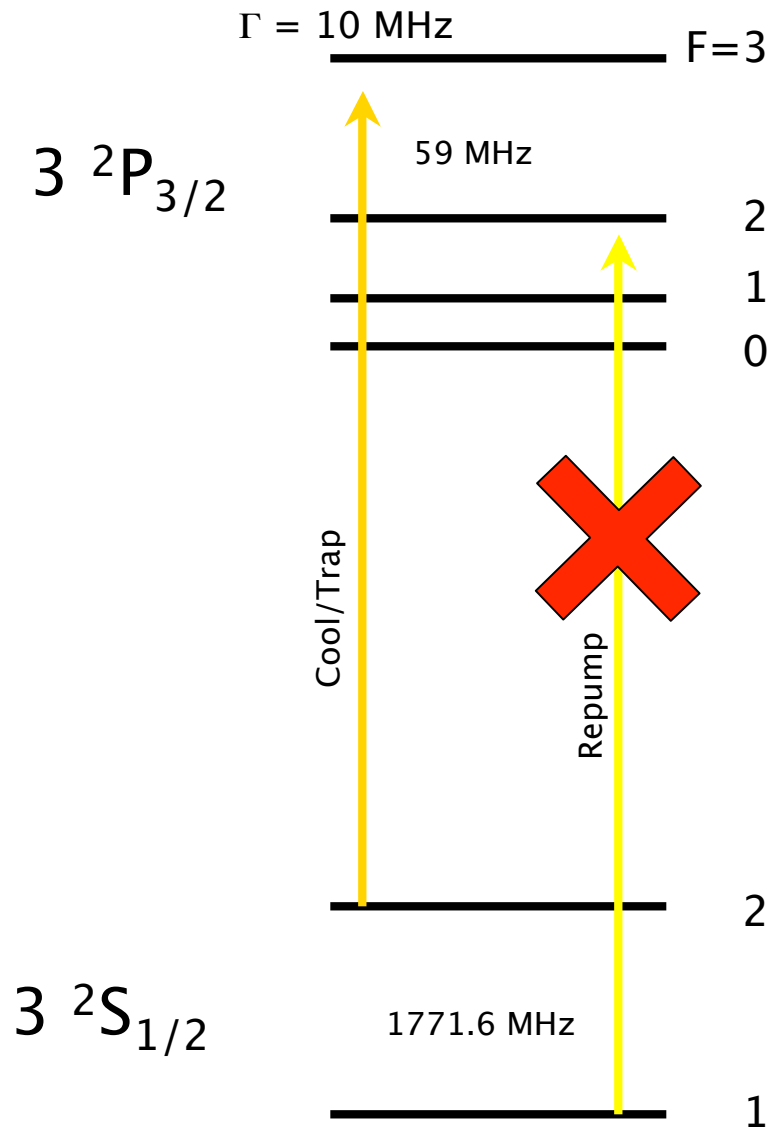
Move the MOT by dragging with magnetic field.

Find the location of the minimum time-of-flight of $^{23}\text{Na}_2^+$ in the horizontal and vertical axis by fitting TOF to a 2-D parabola:



Wire crossing point = (84.74, 86.21)
Simlon (Y,Z)

Reduce molecular formation rate with a Dark Trap



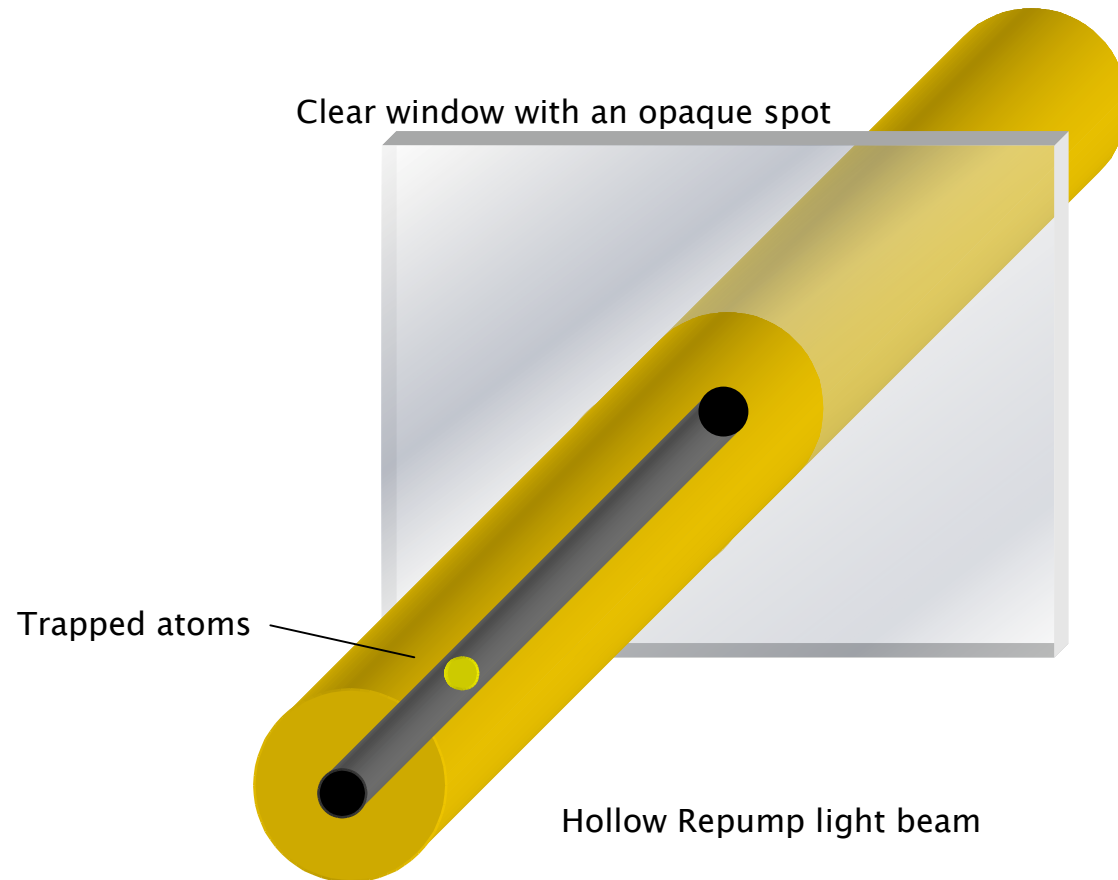
Turn off repump light on trapped atoms

Optically pump to $F_g = 1$

Atoms no longer see resonant light,
Remain in $3\ ^2S_{1/2}$

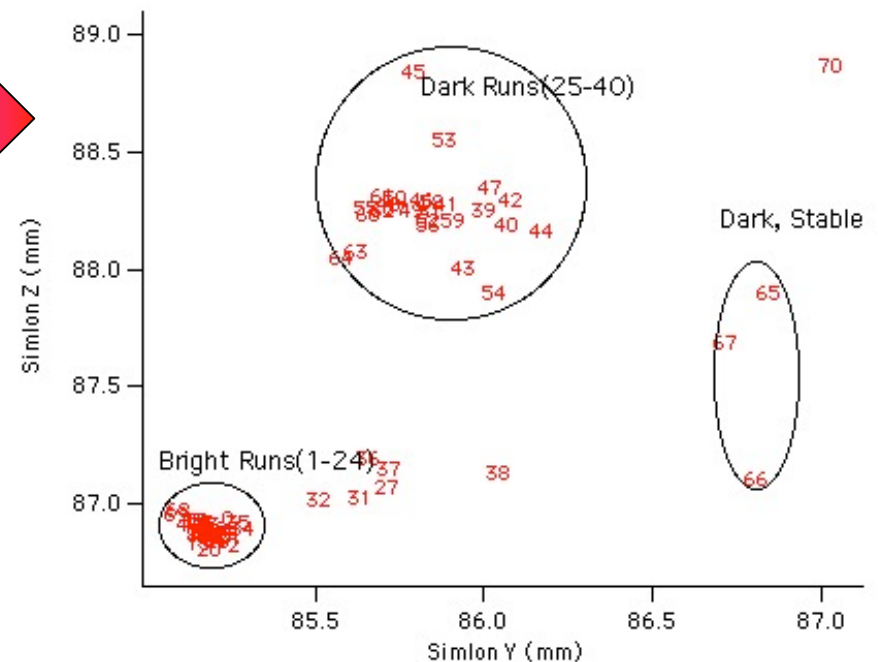
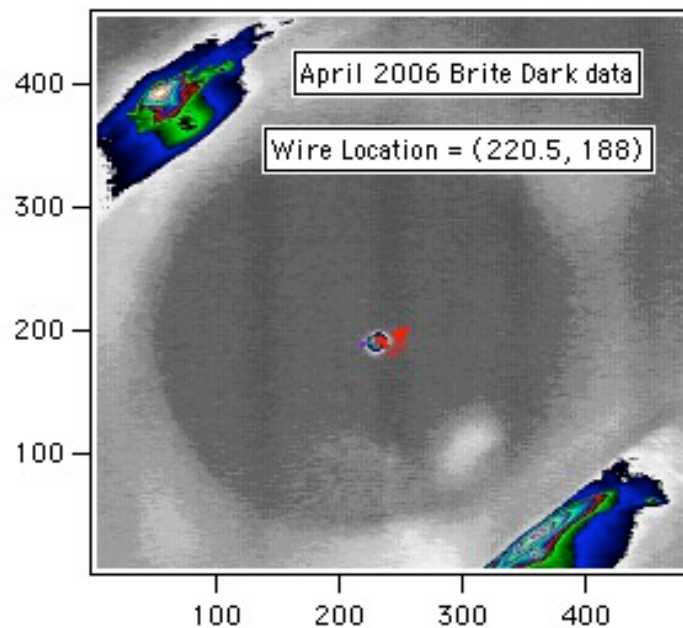
Collisions between ground state atoms
Can't readily form bound molecules.

Dark MOT geometry



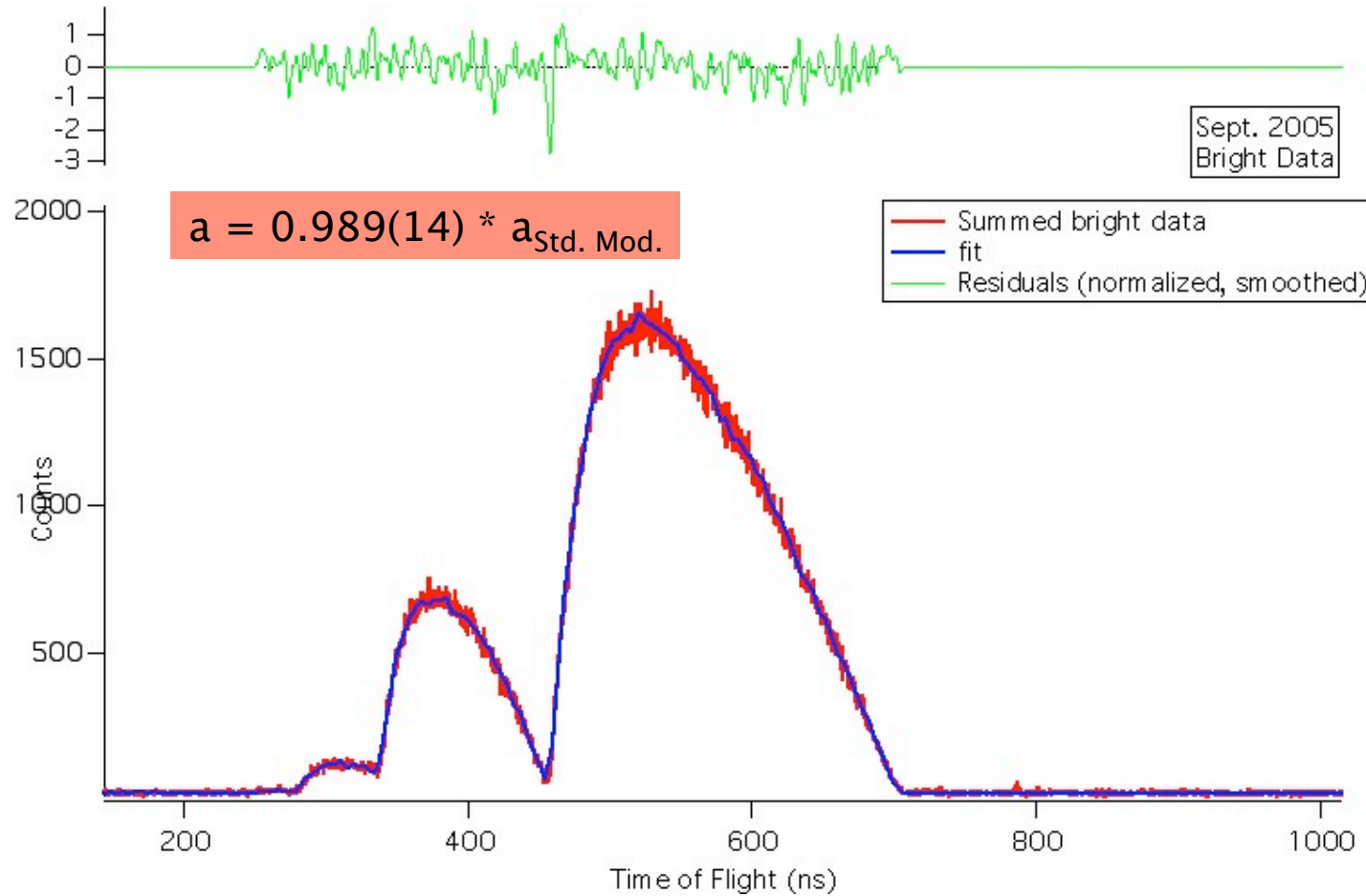
For each data run, determine the trap location

Red dots=location of various runs combined on one representative image

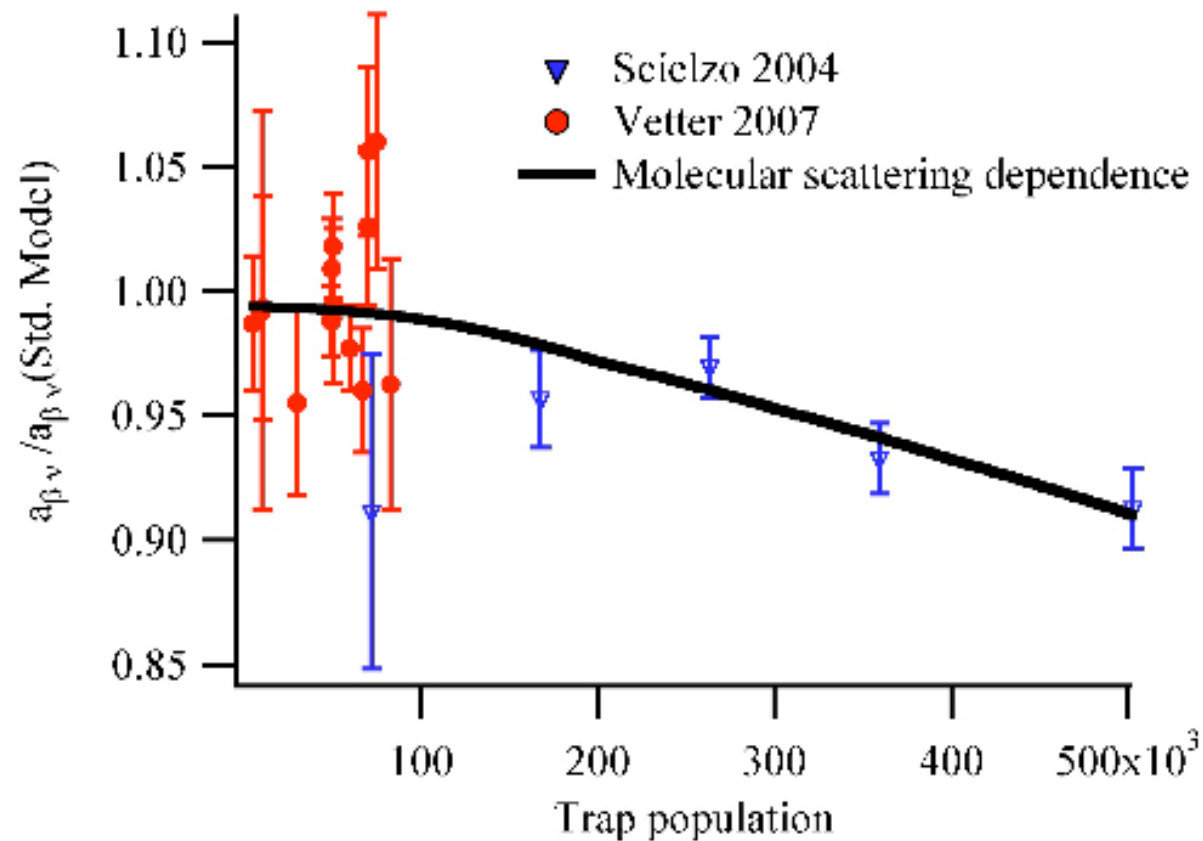


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Sample Data and Fit



Current and previous data



The answer

$$a_{\beta\nu} = 0.9932 \text{ (69)(85) stat sys} * a_{\text{StdMod}}$$

$$a_{\text{StdMod}} = 0.553(2) \text{ Hardy, 2006 + Mass, Q, B.R.}$$

$$a_{\beta\nu} = 0.5502 \text{ (38) (46) (20)}$$

[stat] [sys] [model]

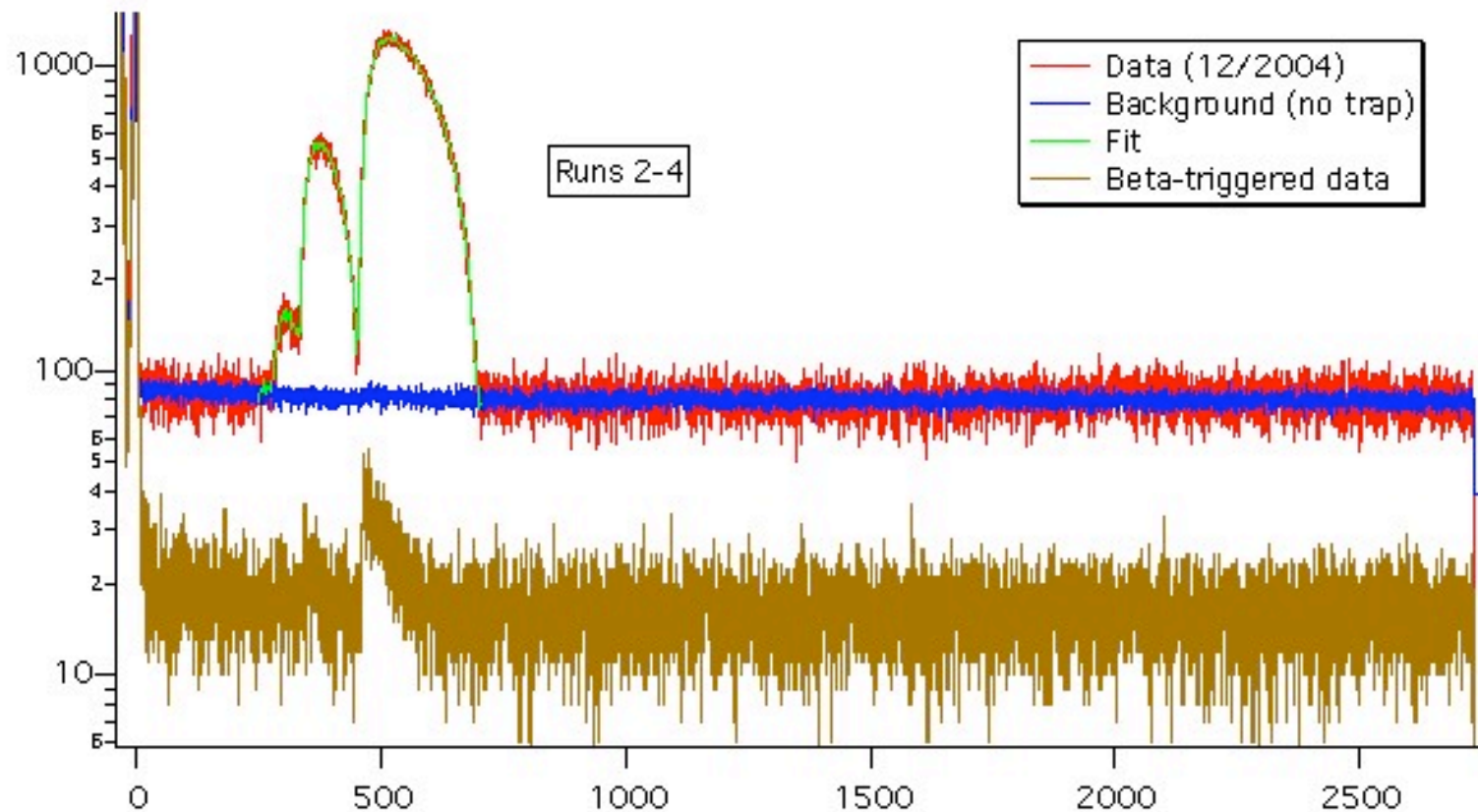
Systematic uncertainties

Most studied via Monte-Carlo comparisons:

generate “clean data” and “dirty data”, then fit “dirty data” using “clean”

Source of Uncertainty	Correction to $a_{\beta\nu}$	Uncertainty in $a_{\beta\nu}$
Beta trigger electron MCP	$\Delta a = -0.5(4)\%$	0.4%
Recoil Ionization	$\Delta a = -0.4(2)\%$ Nick had: $\Delta a = -0.6(3)\%$	0.2%
Polarization/Alignment	0	0.1%
MCP position dependent efficiency	$\Delta a = +0.25(15)\%$	0.15%
Simulation electrode position	0	0.25%
Electric field measurement uncertainty	0	0.1%
MCP diameter uncertainty	$\Delta a = -0.16(16)\%$ (not applied)	0.17%
Trap size	$\Delta a/\Delta(\text{FWHM}) = 0.044$	0.5%
Trap position (X)	$\Delta a/\Delta x = 0.67\%/mm$	0.2%
Trap position (Y, Z)	$\Delta a/\Delta(z,y) = 1.5\%/mm$	+0.3% (one-sided)
Molecular scattering	$\Delta a = -0.05\%$	0.05%
Internal conversion/excited state B.R.	$\Delta a = +0.54\%$	0.13%
Total Uncertainty (Quadrature Sum)	$\Delta a = -0.11\%$	0.85%
Physics (Ft value, branching ratio, recoil order)	$a_{\beta\nu} = 0.554(2)$	0.36%

Beta-triggered event subtraction

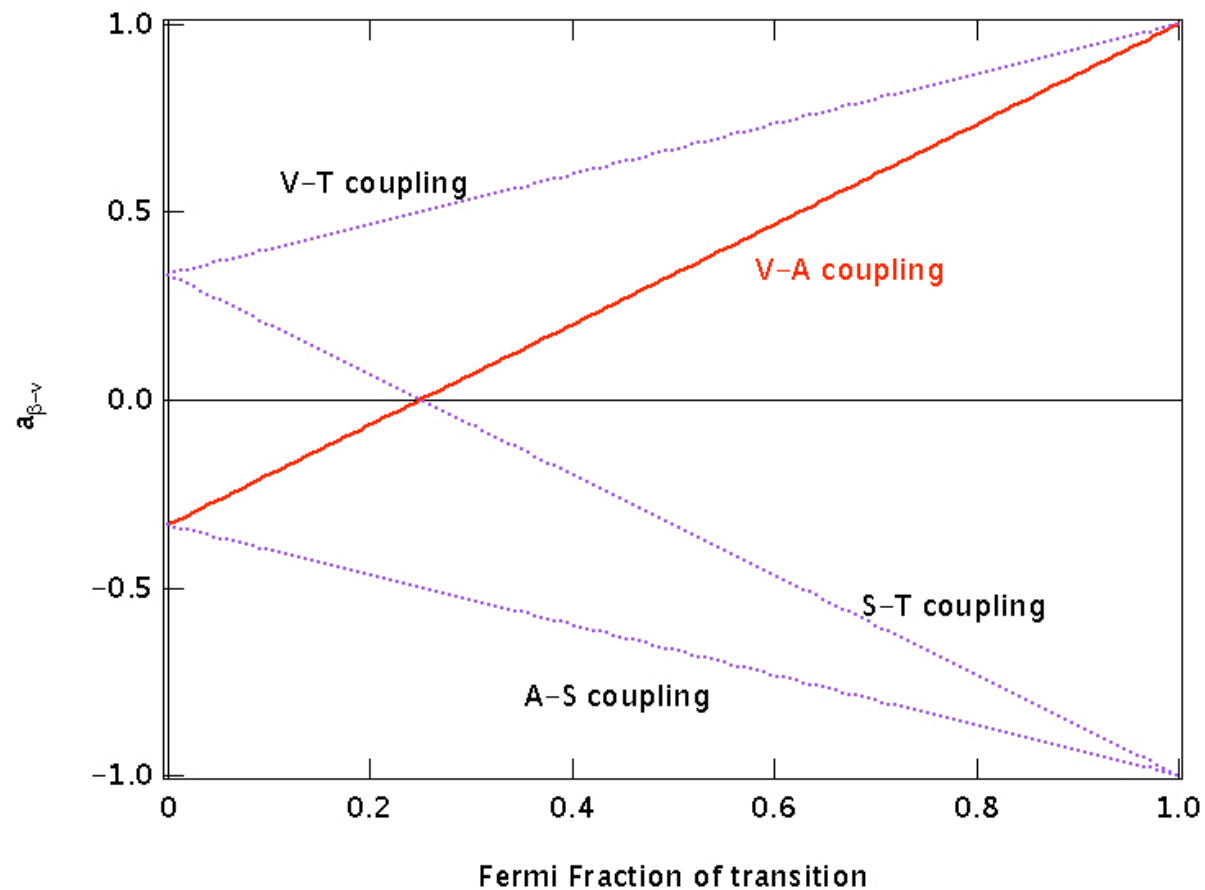


Data shows these events, with reduced efficiency. 3% relative detection efficiency
Unclear how to subtract from data -- need to know absolute detection efficiencies
(multiple electron hits)

Not a large effect, however.

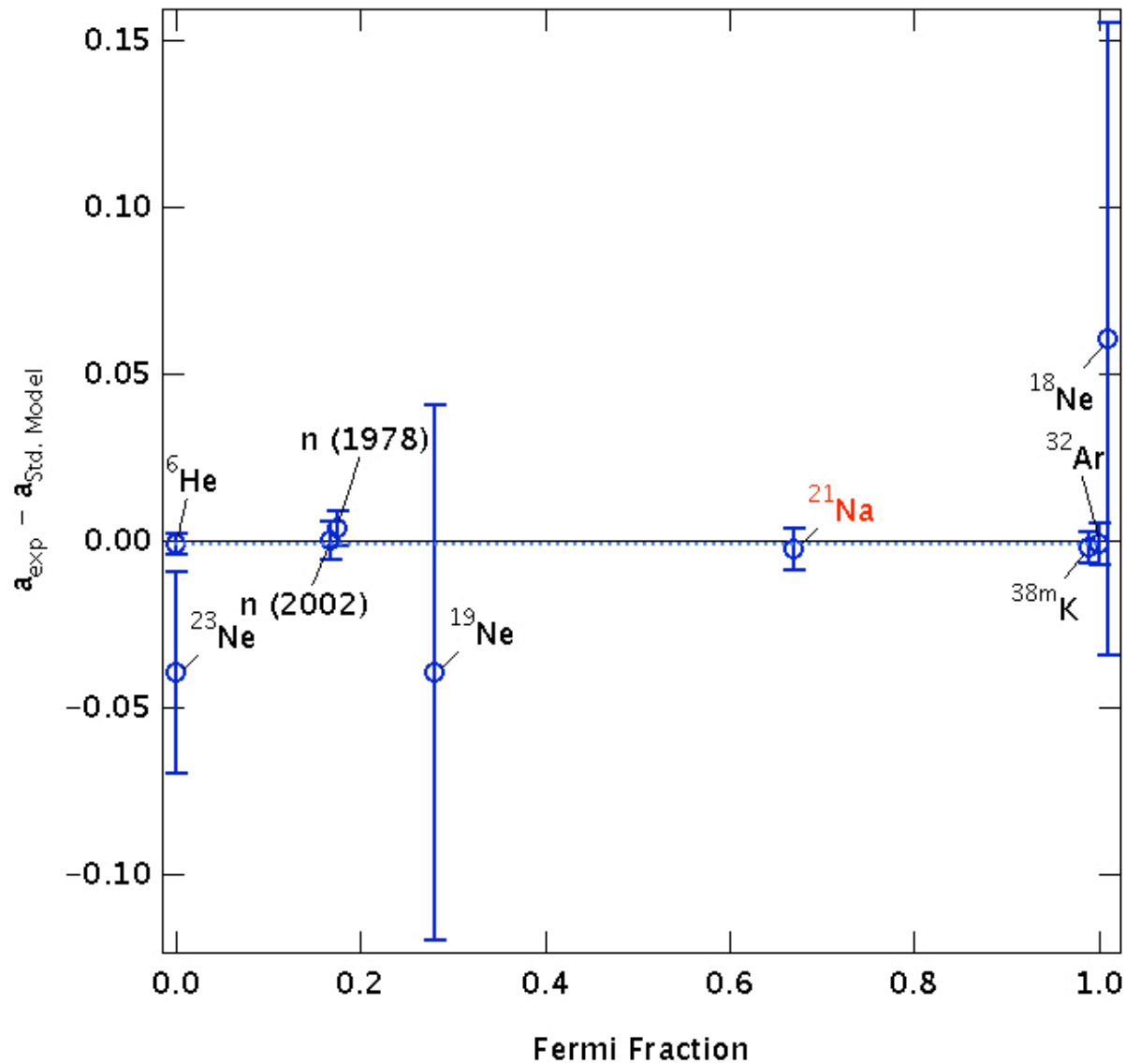
Beta Neutrino Correlation Measurements

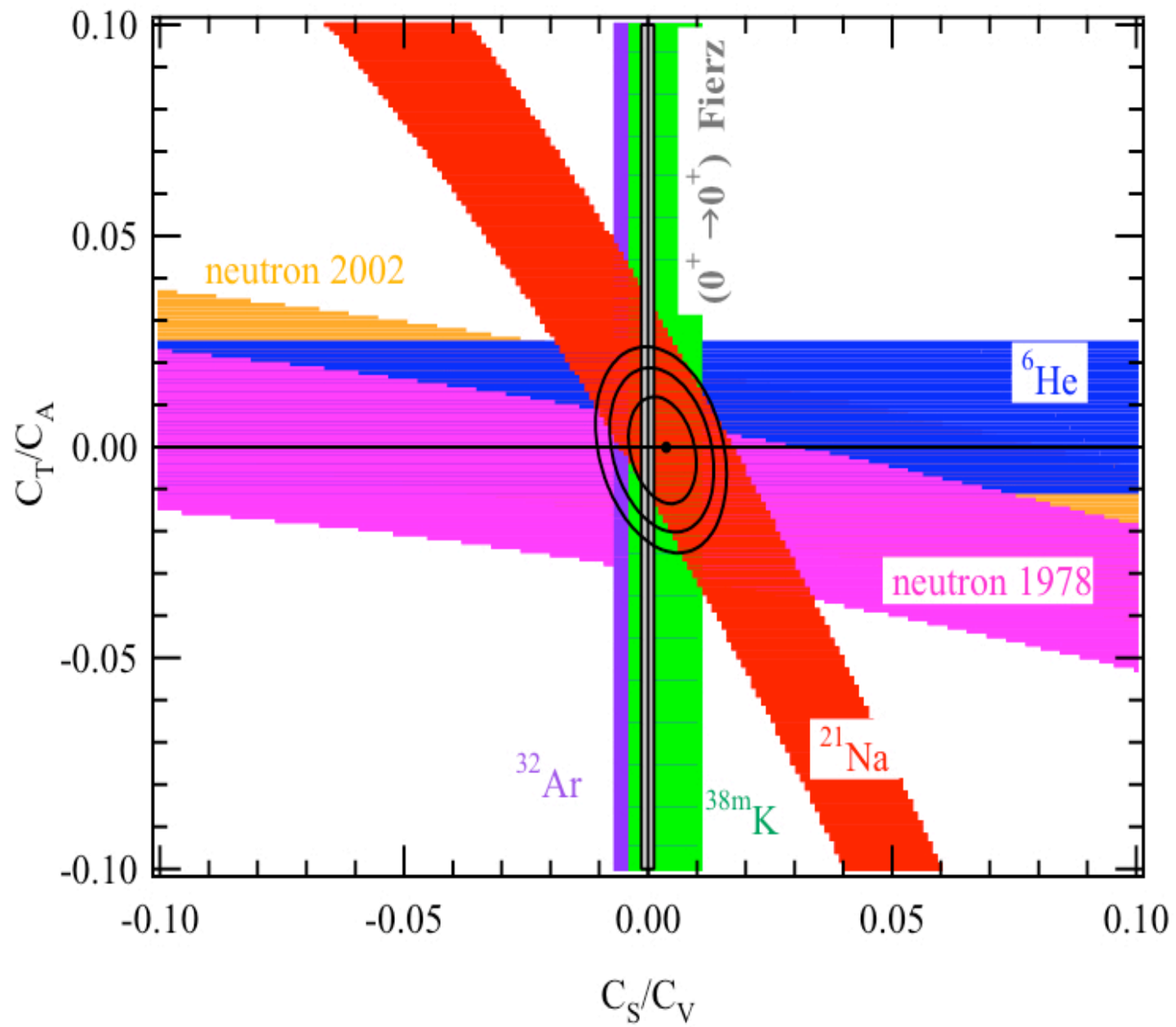
$a_{\beta\nu}$ plotted as a function of the Fermi fraction of the transition:



Beta Neutrino Correlation Measurements

Residuals to $a_{\beta\nu}$ as a function of the Fermi fraction of the transition:





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Sensitivity depends on the Fierz term

$$\frac{d^3\Gamma}{dE_e d\Omega_e d\Omega_\nu} \propto 1 + a_{\beta\nu} \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} + b_{\text{Fierz}} \frac{m_e}{E_e}$$

$$\tilde{a} = \frac{a_{\beta\nu}}{1 + b \frac{m_e}{\langle E_\beta \rangle}}$$

$$b \xi = \pm 2\sqrt{1 - (Z\alpha)^2} \operatorname{Re} [|M_F|^2 C_S C_V^* + |M_{GT}|^2 C_T C_A^*]$$

$$\frac{C_S}{C_V}, \frac{C_T}{C_A} < 0.01$$

Tests for new tree-level physics at $10 M_W$ or $\sim 1 \text{ TeV}$

Implications for BSM physics

Several recent papers, very broad overview of experimental results and sensitivities to new physics

References:

S. Profumo, M.J. Ramsey-Musolf, and S. Tulin, “Supersymmetric Contributions to Weak Decay Correlation Coefficients,” *Phys. Rev. D* **75**, 075017 (2007).

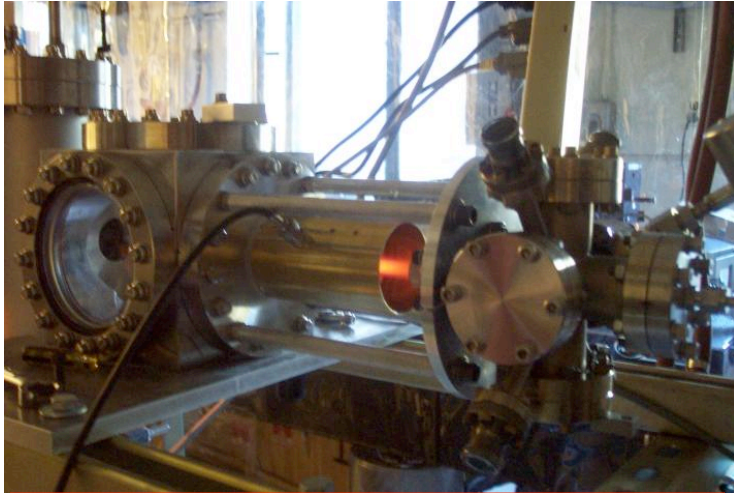
M.J. Ramsey-Musolf and S. Su, “Low Energy Precision Test of Supersymmetry,” hep-ph/0612057.

N. Severijns, M. Beck, and O. Naviliat-Cuncic,
“Tests of the Electroweak Standard Model in nuclear beta decay,”
Rev. Mod. Phys. **78**, 991 (2006).

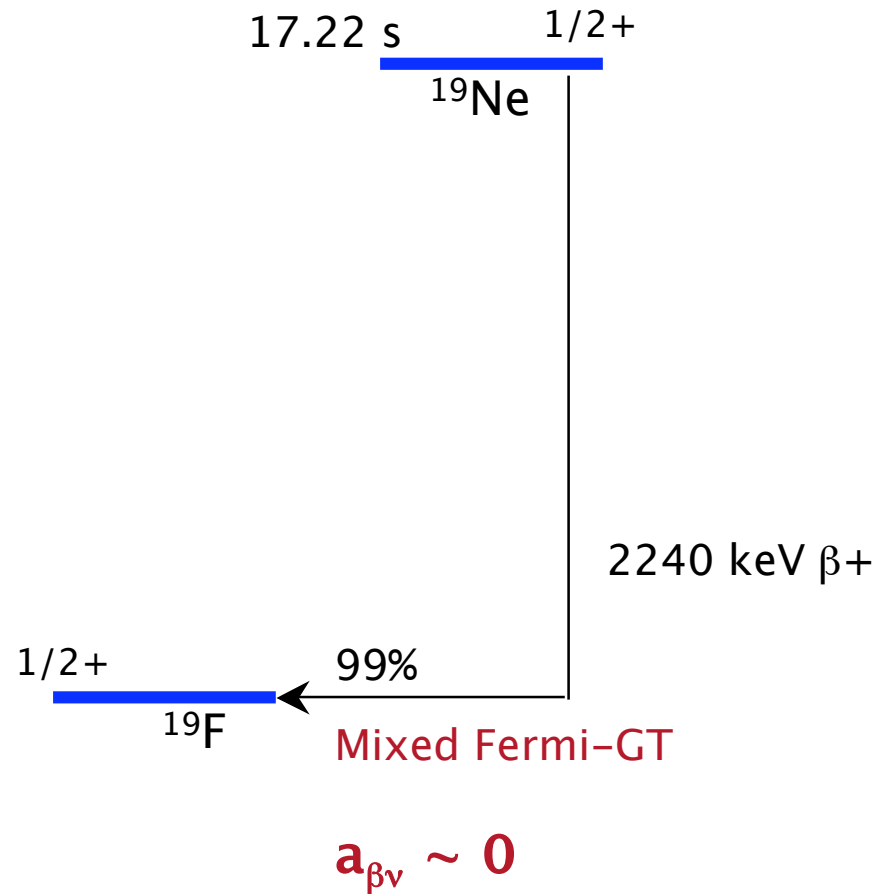
P. Herczeg, “Beta Decay Beyond the Standard Model,” *Prog. Part. Nucl. Phys.* **46**, 413 (2001).

And others... (van Klinken, 1996)

Beta-neutrino correlation in Neon



Metastable neon atomic beam source



Beta-neutrino correlation in Neon

